P PRODUCTION TARGET STUDIES -Numerical Calculations

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INTRODUCTION

The source for antiprotons for the Tevatron I Project should be as "bright" as possible. Since brightness is connected inversely with emittance and since there is a large angular emittance of the antiprotons due to their production transverse momentum, it is desirable to make the source have the smallest practical spatial size. Two basic considerations limit the spatial dimensions of the source, finite density of real target materials and multiple coulomb scattering in the target itself.

The first effect introduces a "depth of field" optical abberation and the second increases both the angular spread of the source and the effective source size. In practical schemes, the instantaneous and average energy deposition in the target by the proton beam which makes the antiprotons also introduces a "depth of field" optical abberation and the second increases both the angular spread of the source and the effective source size. In practical schemes, the instantaneous and average energy disposition in the target by the proton beam which makes the antiprotons also introduce serious (and to some extend separable) constraints on the problem. Finally, the proton beam incident has a finite emittance of its own which has an impact on the result separable to some extent from the target heating problems. Many people have made estimates and analytic formulations of various combinations of the above factors and drawn conclusions of seemingly contradictory nature. Since a complete numerical calculation of the p emittance, including all of the effects in an accurate manner, is possible, it seemed useful to do this calculation and summarize the results for use by optical transport programs concerned with capturing \bar{p} 's for cooling and storage.

In this study, the target density will be assumed to remain constant since exploding targets will be very unlikely in a practical scheme. It is further clear that the overall production cross section for \bar{p} 's factors out of the problem and is therefore not addressed explicitly here. The production transverse momentum, on the other hand is a crucial factor, and a model for this is introduced. Likewise, chromatic effects due to finite momentum spread of the \bar{p} 's is not explicitly carried along, but its effect is presented in one study here.

INPUT MODELS and DATA

Several elements are assumed in making a study of practical use. Let us

- a) Proton beam: The incident proton beam is assumed to have a symmetric x and y gaussian emittance in x and x' given by
 - ε_{v} = ε_{x} = π $\sigma_{x}\sigma_{x}$ = 0.05 π mm mr. (1 standard deviation)

It is also assumed that the beam size and divergence can be varied within the constraint that phase space(emittance) is conserved.

- b) <u>Target</u>: The target is assumed to be tungsten(but if rhenium is used, its parameters are virtually identical) of transverse size large compared to the beam spot. In particular, all particles are assigned multiple scattering(and nuclear attenuation) as though they cannot escape from the sides of the target.
- c) Production transverse momentum the antiprotons are assumed to have a production transverse momentum wrt the proton direction of:

$$\frac{dN}{dP_{\perp}^{2}} = C_{0} e^{-(P_{\perp X}^{2} + P_{\perp y}^{2})/2\sigma_{\perp}^{2}}, \sigma_{\perp} = 0.40(\text{GeV/c}).$$

- d) Target length: Target lengths of 0.0, 2.0, 4.0, 6.0, 8.0 cm are considered.
- e) Target efficiency: The yield per incident proton per \bar{p} produced is lowered by attenuation in the target of the incident protons and of the produced \bar{p} 's. This efficiency factor F_{tgt} is given by:

$$F_{tgt}(\lambda_p, \lambda_{\bar{p}}, \ell) = \frac{e^{-\ell/\lambda_p} - \ell/\lambda_{\bar{p}}}{(\lambda_p/\lambda_{\bar{p}} - 1)}$$
,

where; ℓ = target length(cm)

 λ_p = absorbtion length for high energy(nom 80 GeV) protons(cm)

 $\lambda_{\overline{p}}$ = collision length for low energy(nom 5 GeV) antiprotons(cm).

note:

if
$$\lambda_p = \lambda_{\overline{p}}$$
, then,
 $F_{tgt}(\lambda, \ell) = \frac{e^{-\ell/\lambda}}{\lambda/\ell}$.

We choose $\lambda_p \stackrel{\text{\tiny if}}{=} 10$ cm, 5 $^{\varsigma}$ $\lambda_{\overline{p}}^{-}$ $^{\varsigma}$ 10 cm .

f) <u>Multiple scattering</u>: The projected multiple scattering in a "thick" target is a correlated two dimensional gaussian in position and slope, and is given by:

$$\frac{d^{2}P(\rho, x', t)}{dxdx'} = C_{0}e^{-B(x'^{2}-3x'\rho + 3\rho^{2})}$$

where;

x' = change in slope(angle) due to multiple scattering in a distance t

ρ = (x/t)= change in transverse position x from the unscattered orbit by traversing a thickness t of multiple scattering material

$$B = 4 \left(\frac{\beta P}{E_S} \right) \left(\frac{x_0}{c} \right)$$

 $E_S = .021 \text{ (GeV)}$

P,E = particle momentum(GeV/c), total energy(GeV)

 $\beta = P_*E = Lorentz velocity parameter$

 x_0 = radiation length of scatterer(cm)

 $x_0 = 0.35$ cm (tungsten), 155.0cm(lithium)

t = thickness of scatterer (cm)

 $C_0 = \frac{2\sqrt{3}!}{\pi} \frac{x_0(\beta P)^2}{t^2(E_S^2)} = \text{normalization constant.}$

When applying multiple scattering to an optical system, the position/angle correlation must be explicitly acknowledged or erroneous results may be obtained. The spatial multiple scattering distribution accurately factors into the product of the projected multiple scattering in the two perpendicular directions for small angle trajectories(the case at hand).

g) <u>Lithium lens</u>: The antiproton beam emerges from the production target at 5.35 GeV/c momentum and with relatively large transverse momentum in the current Tevatron I plan. To capture a large fraction of these p̄'s, a monopole lens made from lithium carrying a large pulsed current is proposed. The multiple scattering and nuclear attenuation in the lens is of interest, so its effects have been included. We describe the lens by:

L = length = 10.0 cm

G = magnetic gradient = 100.0 KG/cm

With these elements, we can monte-carlo the emittance of the \bar{p} 's. The method is to choose 10,000 rays from a gaussian proton beam phase space in x and x', assume a \bar{p} produced uniformly along a target of length T with a projected transverse momentum chosen according to an gaussian distribution in P_{\perp} as in c), and propagate this ray with multiple scattering through the remaining target, across a 20.0 cm drift space(measured from the <u>center</u> of the production target) and on through the lithium lens(also with multiple scattering). From this point, the ray can be propagated forward to the focus of the lithium lens or backward (without scattering) to the center of the production target. In both these locations, the <u>thin target</u> phase ellipse in(x,x') is erect and the one dimensional distributions for a thin target slice are simple.

The actual distributions are <u>not</u> simple for targets of length several centimeters due to depth of field distortions. That, of course, is the whole reason for doing the problem numerically!

RESULTS

The most direct way to present the results of the monte-carlo calculation is in the form of a density plot of numbers of \bar{p} 's in small bins of x and x', the transverse position and slope of \bar{p} rays from the production target. As noted above, all rays are presented as they reproject to the center of the production target. In some cases, the detailed distribution at the "hot spot" near x = 0.0, x' = 0.0 is presented in a second plot with finer bin size.

In each of the plots in Figures 2-19, the numbers shown in each bin in the (x, x') plot represent those particles out of a sample of 10,000 produced \bar{p} 's that ended up in a particular piece of phase space. Note that this is a <u>two</u> dimensional study in <u>one</u> transverse coordinate, hence the yields into a three dimensional solid angle go as this projected result squared only for small solid angles. Since real beam transports will find it hard to subtend large parts of the phase space, the projected results are nevertheless useful.

The plots labeled "Correlation Plot No.3" have bins of 0.04 cm in x and 10.0 milliradians in x'. The border around the central 400 bins

gives the overflow events that do not fit into the scale of the plot. The bottom row and rightmost column give the sums of their respective columns and rows, respectively. The plots labeled "Correlation Plot 4" give the same results in a magnified central region with x bins of 0.02 cm and x' bins of 4.0 mr. The parameters of beam, target and Li lens are included on the plot. Multiple scattering in the target and lens are included unless noted otherwise.

Figure 1 shows the geometry of the problem and Table I summarizes the cases plotted in Figures 2-17.

TABLE I

Correlation Plots for p Production in Tungsten Targets

Figure Number	Beam RMS Width (mm)	Beam RMS Slope (mr)	Target Length(cm)	Target Multiple Scattering	Li Lens Multiple Scattering	Comments
2	0.300	0.167	0.0	YES	YES	Target length study
3	H	п	0.0	п	п	with all multiple scattering and
4	n	n ·	2.0	н	п	"large" incident
5	н	tt	4.0	н	н	proton beam
6	11	ŧ)	4.0	B	n n	
7	11	п .	6.0	II.	11	
8	н	11	8.0	II	я	
. 9	11	tt.	8.0	11	11	
10	0.075	0.067	4.0	.		"Small" beam study
11	11	H	4.0	II	11	with full multiple scattering
12	0.0	0.0	0.0	NO	ЙО	Phase space for point p source
13	11	n	0.0	11	II	
14	п	II	4.0	u	11	Finite target length effect
15	· u	11	4.0	н .	. 11	
16	11	<u>ii</u>	0.0	11	YES	Effect of multiple
17	n	н	0.0	н	YES	scattering in LiLens
18	0.300	0.167	4.0	YES	YES	$\Delta P/P = \pm 2\%$ added to
19	n	H	4.0	н	- n	show chromatic behavior

INTERPRETATION

The raw phase space density data contained in the plots of Table I can be presented in a number of useful ways. Before combining the effects of beam attenuation in the(finite) target with attenuation of the produced \bar{p} 's in the same target, it is useful to study the effects of proton beam radius and target length normalized to the <u>same</u> number of \bar{p} 's emerging from the target. Later, the various attenuation and production factors can be multiplied by the densities to get the actual production members. In all cases we study, the full multiple scattering of the target and lithium lens is included.

In Figure 20, the phase space density of the "hot spot" at x = 0.0, x' = 0.0 is plotted versus target length for several incident proton beam radii. The central brightness varies weakly with target length over the range 0-8.0 cm. Likewise, there is a gentle dependence on proton beam radius for radii in the range 0-0.30 mm.

This latter result is seen more clearly in Figure 21 where the central density is plotted versus rms beam radius for several target lengths. Note that the beam phase space density is held constant for each beam radius as is necessary for any real beam focussing system. The phase space density of the incident proton beam is assumed to be 0.05 π mm mr as noted above. The central phase space density is a concept of limited usefulness. It is primarily valuable for small \bar{p} phase space acceptances, in which circumstances the yield into an acceptance $\epsilon_{\chi}\epsilon_{V}$ is given by:

$$Y \stackrel{\cong}{=} \frac{d^{2}N}{dxdx}, \int_{0}^{1} \varepsilon_{x} \varepsilon_{y}$$

$$= \int_{-y}^{+y} \int_{-y}^{+y} \int_{-x}^{+x} \int_{-x}^{+x} dxdx'dydy' \frac{d^{4}N}{dxdx'dydy'}.$$

For finite solid angle acceptance of \bar{p} 's, the numerical integral over discrete regions of phase space is more useful. In the case of Figures 2-19, this integral is proportional to the sum over the desired ranges of x and x'. In the limit where multiple Coulomb scattering and \bar{p} transverse momentum are both gaussian in the transverse spatial directions(a good approximation for all cases of practical interest), the full four dimensional phase acceptance goes as the product of the(x,x') and (y, y') projections. This being the case, each count in a correlation plot bin represents 1.0 x 10^{-4} of the \bar{p} 's produced in the target.

Figure 22 shows the fractional yields as a function of ϵ_X for symmetric (x, x'), (y,y') values of the \bar{p} emittance. The curves are parametric in target length and assume a proton beam of 0.30 mm x 0.167 mr. Full multiple scattering is included; the beam and \bar{p} attenuation factors are not i ded at this point. The curves verify the unremarkable fact that large acceptance means efficient collection while small acceptance means small yield. This is not news, but it is useful to see the dramatic increase in \bar{p} collection as ϵ_X varies from 5π to say 50π mm mr.

Note carefully that the <u>emittance</u> integral is over <u>rectangular</u> regions in (x, x') phase space rather than the conventional elliptical ones. When comparing these emittances to an <u>acceptance</u> that is <u>elliptical</u> in (x, x'), but which is nevertheless limited by discrete angular limits, a factor of $(\pi/4)$ may be needed to obtain the elliptical <u>admittance</u> from the rectangular <u>emittance</u> for fixed angular boundaries.

The effect of finite target length on beam and \bar{p} attenuation is shown in Figure 23. The best model for proton beam and \bar{p} attenuation probably lies between the two extreme cases shown. If elastic scattering of protons is ignored(a good assumption) then λ = 10 cm. If elastic scattering of \bar{p} 's is ignored(a questionable, but not unreasonable assumption), then $\lambda_{\bar{p}}$ = 10 cm. If elastically scattered p's are assumed lost, $\lambda_{\bar{p}}$ = 5 cm. Real life probably lies nearer the $\lambda_{\bar{p}}$ = 10 cm limit.

In Figure 24, the target attenuation factor from Figure 23 multiplies the fractional yields of Figure 22 to produce the actual target length dependent yields in tungsten. To obtain the absolute yield of \bar{p} 's, it is only necessary to multiply the ordinate shown by the fraction of inelastic proton collisions that produce one or more antiprotons in the accepted \bar{p} momentum interval. If the momentum acceptance exceeds a few precent, integration of the production over the momentum interval may be necessary.

The \bar{p} yield is finally the important quantity of interest. It can be given by the equation:

$$N_{\bar{p}} = N_{p} t f_{TGT} \Delta \sigma = 2\pi N_{p} \left(\frac{E d^{3} \sigma}{dP_{3}} \right) \sigma_{L}^{2} \Delta P_{L} \Delta A(t, \epsilon)$$

where; $N_{\bar{p}}$ = no. of \bar{p} 's produced from target t with mean total energy E, in the momentum interval ΔP_{ii} , captured in the symmetric emittance

$$\varepsilon_{x} = \varepsilon_{y} = \varepsilon$$

 $N_{\rm p}$ = no. of beam protons incident on the target

$$t = N_{0} \rho \ell$$

 $N_0 = 6.02 \times 10^{23}$ (Avagadro's Number)

 ρ = density of tungsten (gm/cm³)

 ℓ = target length (cm)

 f_{TGT} = finite target length factor (see Figure 23)

$$\Delta\sigma = \int_{0}^{P_{\perp}^{2}} \max_{\mathbf{max}} \int_{P_{\parallel}}^{P_{\parallel}} \max_{\mathbf{min}} \frac{d^{2}\sigma}{dP_{\perp}^{2}dP_{\parallel}} dP_{\perp}^{2}dP_{\parallel}$$
 (cm²)

 $E = \bar{p}$ total energy (GeV)

 ΔP_{H} = accepted \bar{p} momentum bite(GeV/c)

 $\sigma_{\underline{I}}^2$ = variance for the assumed \bar{p} production transverse momentum (assumed gaussian), (GeV/c)²

$$\left(\frac{\text{Ed}^3\sigma}{\text{dP}^3}\right)$$
 = invariant per nucleon p production cross section evaluated at P_⊥ = 0, E (cm²/GeV²)

 $A(t,\epsilon)$ = all the geometrical and target length effects. This is the factor that must be computed numerically. Its value varies from 0 to 1.0; it is dimensionless.

Several of the numbers needed for numerically calculating the \bar{p} yields are uncertain. The values for σ_L^2 and $\left(E\frac{d^3\sigma}{dP^3}\right)$ are known only approximately.

Nevertheless, by assuming specific values, we can demonstrate \bar{p} yields explicitly. The values we choose are given by:

$$N_p = 10^{13}$$
 $\sigma_L^2 = 0.16 \text{ (GeV/c)}^2$
 $\left(\frac{\text{Ed}^3\sigma}{\text{dP}^3}\right) = 1.0 \frac{\text{mb}}{\text{GeV}^2} = \frac{1.0 \times 10^{-27} \text{cm}^2}{\text{GeV}^2} \text{ @ E = 5.4 GeV}$
 $\Delta P_{\text{cs}} / \text{E= 0.04}$

With these values, the actual \bar{p} yields are shown on the right hand scale in Figure 24. From the graph, the emittance and target length dependences are quite clear.

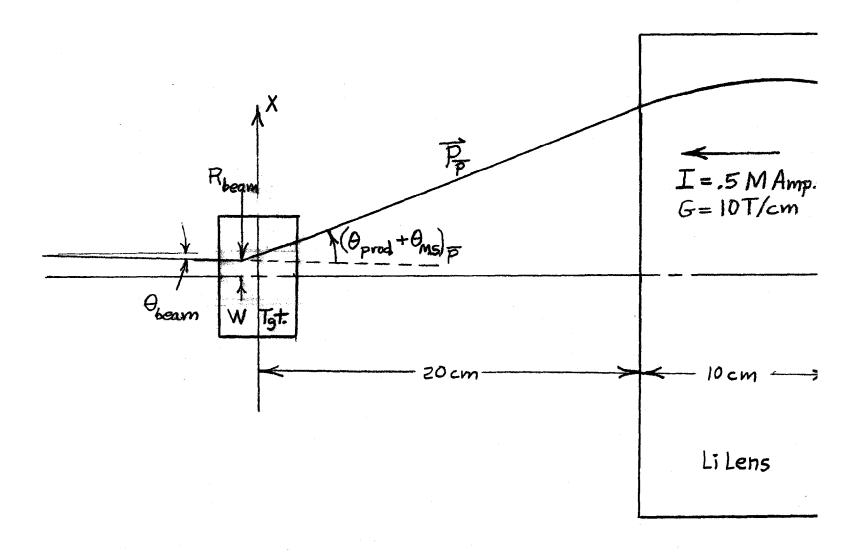


Fig 1 P Target Geometry

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Fig 6

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Fig9

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